

# The semi-inclusive approach to measuring $x_J^\gamma$

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# Factorizing jet signal and uncorrelated background: experiment

## Heavy ion collisions

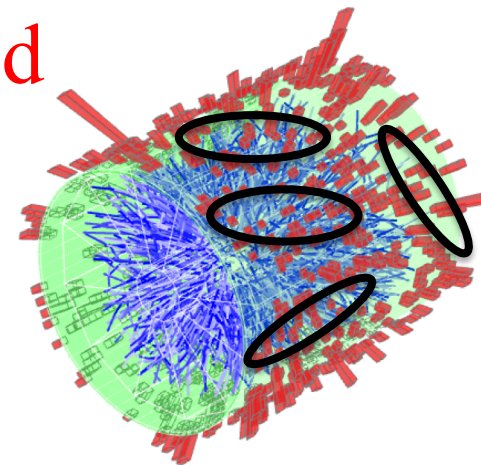
- Multiple hard processes generating jets
- Copious soft hadron production
- Many such processes contribute to each reconstructed jet candidate

## Define the signal jet population

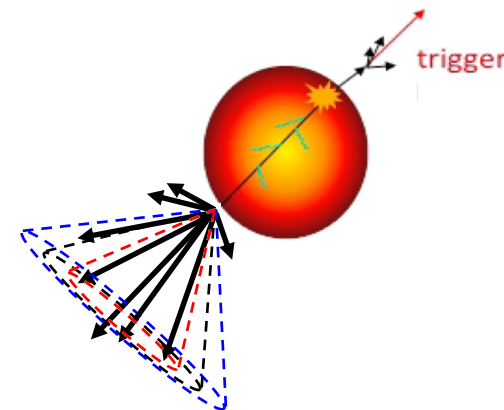
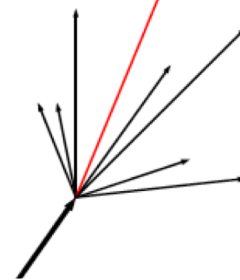
- require linear response in bkgd environment  
→ must be rare  $\ll 1/\text{event}$
- can then correct for  $p_T$ -smearing via unfolding

## Discriminate hard jet candidates from bkgd

- Inclusive distributions:
  - hard core (high  $p_T$  hadron or cluster) - bias?
  - or high jet  $p_T$  relative to bkgd fluctuations
- Coincidence distributions:
  - Trigger on a hard process (hadron, jet, gamma, Z)
  - Measure correlated energy-momentum distribution recoiling from the trigger
- Jet substructure: grooming – bias?

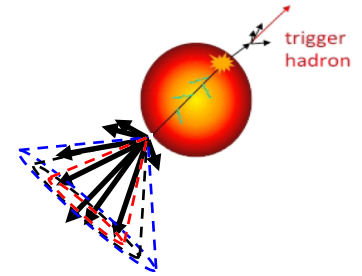


$$p_{T,lead} > p_{T,lead}^{min}$$



# Fully corrected spectrum: illustration of procedure

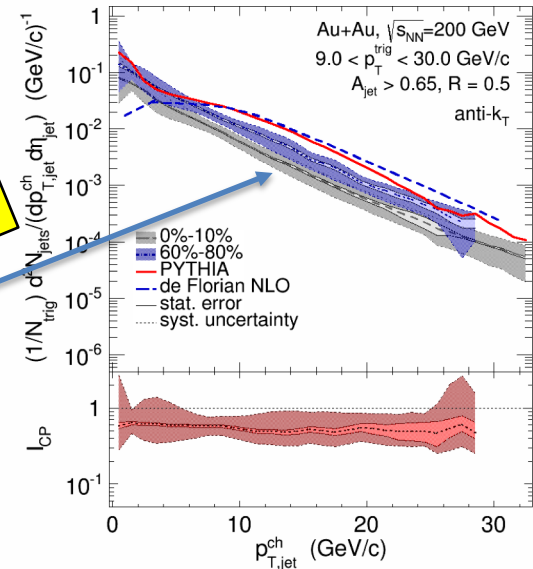
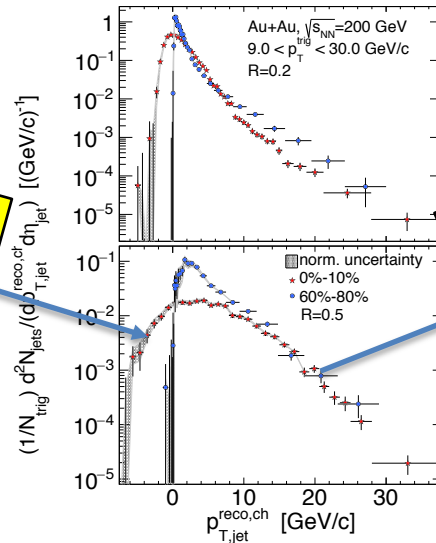
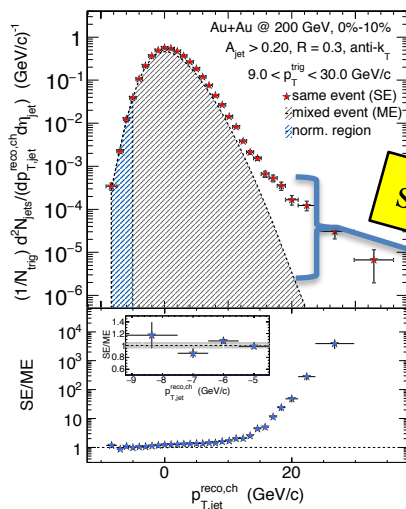
- Cannot know with certainty whether any given hadron (or hadron cluster) arises from the same high- $Q^2$  process as the trigger
- “uncorrelated background” strictly has meaning only for ensemble-averaged distributions



Two distinct steps for correcting bkgd:

Correct for uncorrelated yield (“vertical correction”)

Correct for  $p_T$ -smearing of true jet signal (“horizontal correction”)



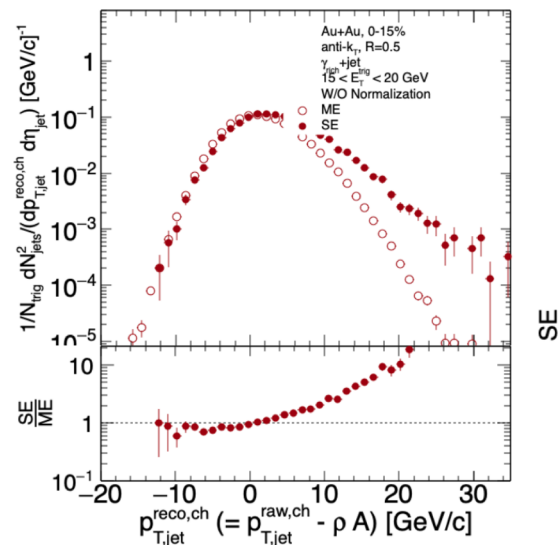
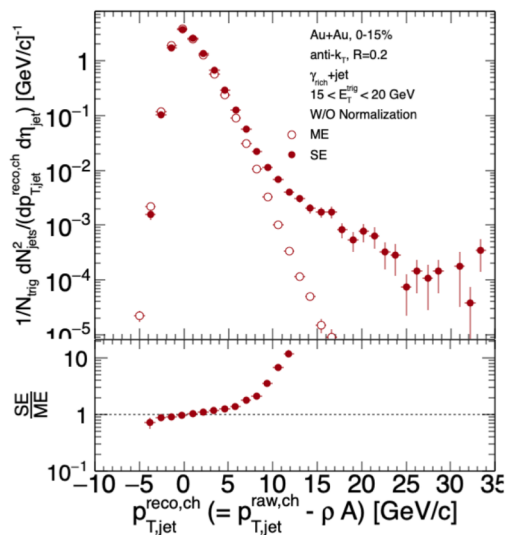
Procedure is driven by the fact that in experiment we cannot know the precise uncorrelated distribution in each event

# Semi-incl gamma+jet: current status

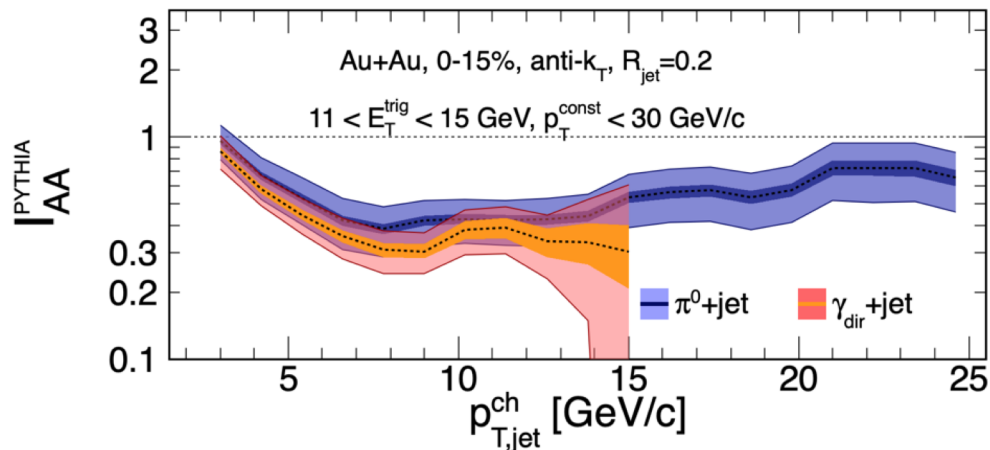
Run 14 data, 13 nb<sup>-1</sup>

$$\frac{1}{N_{trig}^\gamma} \frac{dN_{jet}^{recoil}}{dp_{T,jet}} = \frac{1}{\sigma^{AA \rightarrow \gamma+X}} \frac{d\sigma^{AA \rightarrow \gamma+jet+X}}{dp_{T,jet}}$$

Raw data/ME comparison  
(ME not renormalized)



I<sub>AA</sub> for corrected spectra  
(Nihar's talk at Collab mtg)



# Semi-inclusive $x_J^\gamma$

Current measurement:

$$\frac{1}{N_{trig}^\gamma} \frac{dN_{jet}^{recoil}}{dp_{T,jet}} = \frac{1}{\sigma^{AA \rightarrow \gamma + X}} \frac{d\sigma^{AA \rightarrow \gamma + jet + X}}{dp_{T,jet}}$$

Scale the  $p_T$  of each recoil jet candidate by  $p_{T,\gamma}$   
( $p_{T,\gamma}$  is measured well for each trigger)

$$x_J^\gamma = \frac{p_{T,jet}}{p_{T,\gamma}}$$

Semi-incl distribution of  $x_J^\gamma$

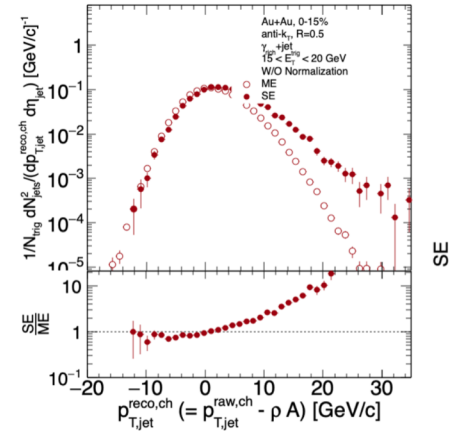
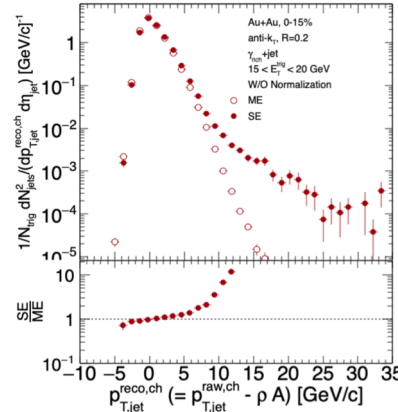
$$\frac{1}{N_{trig}^\gamma} \frac{dN_{jet}^{recoil}}{dx_J^\gamma} = \frac{1}{\sigma^{AA \rightarrow \gamma + X}} \frac{d\sigma^{AA \rightarrow \gamma + jet + X}}{dx_J^\gamma}$$

In practice: measure recoil jet distributions for  $\gamma$ -rich and  $\pi^0$  triggers and disentangle  
This is a “technical detail” ;-), above procedure still applies

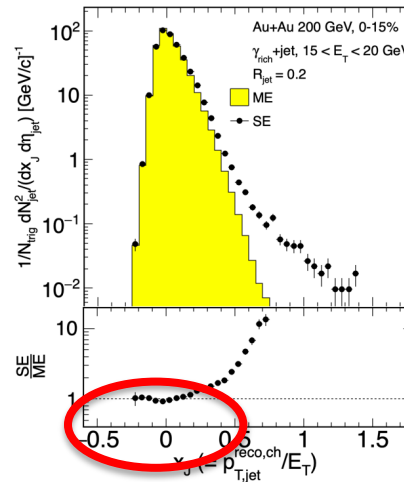
# $\gamma^{\text{rich}}$ data in central Au+Au: SE + ME

*ME not renormalized in these figures*

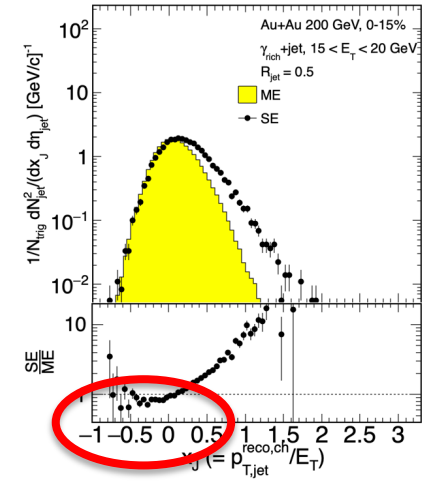
$$\frac{1}{N_{\text{trig}}^{\gamma}} \frac{dN_{\text{jet}}^{\text{recoil}}}{dp_{T,\text{jet}}} = \frac{1}{\sigma^{AA \rightarrow \gamma+X}} \frac{d\sigma^{AA \rightarrow \gamma+jet+X}}{dp_{T,\text{jet}}}$$



$$\frac{1}{N_{\text{trig}}^{\gamma}} \frac{dN_{\text{jet}}^{\text{recoil}}}{dx_J^{\gamma}} = \frac{1}{\sigma^{AA \rightarrow \gamma+X}} \frac{d\sigma^{AA \rightarrow \gamma+jet+X}}{dx_J^{\gamma}}$$

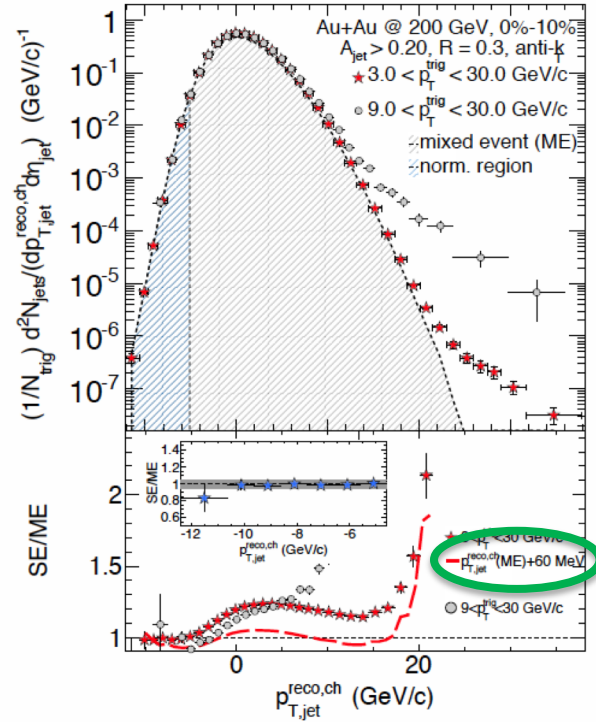
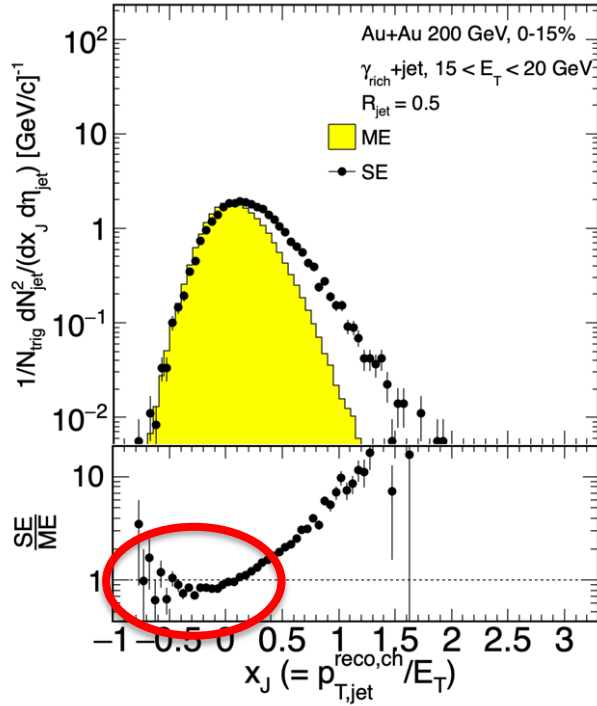


good



see next slide

# h+jet paper Fig. 9



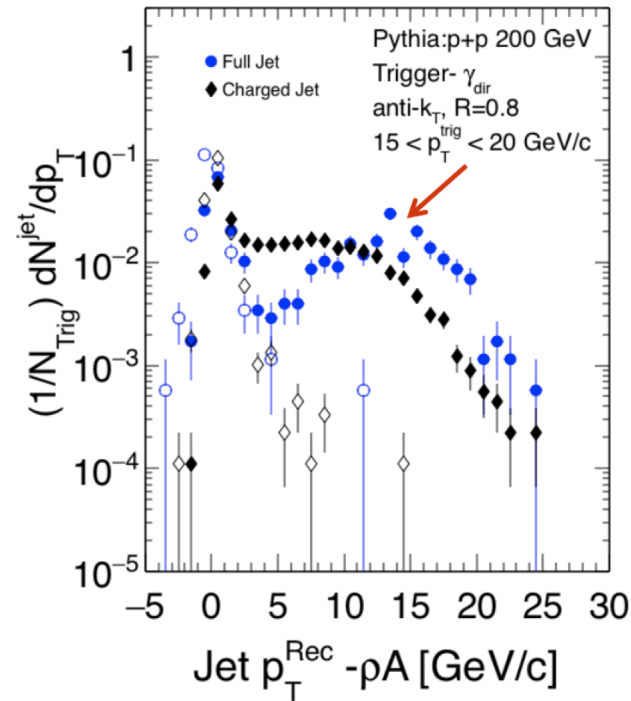
$$\rho = \text{median} \left\{ \frac{p_{T,jet}^{raw,i}}{A_{jet}^i} \right\}$$

Jet populations are different in SE and ME

- small misalignment of  $\rho$  to be expected
- correctable with high precision

Overall: ME correction approach works well for  $x_J \gamma$ ; needs a bit of tuning (as expected)

# Main open issue: full jet measurements w/ BEMC



Charged vs. full recoil jet: significant difference in resolution

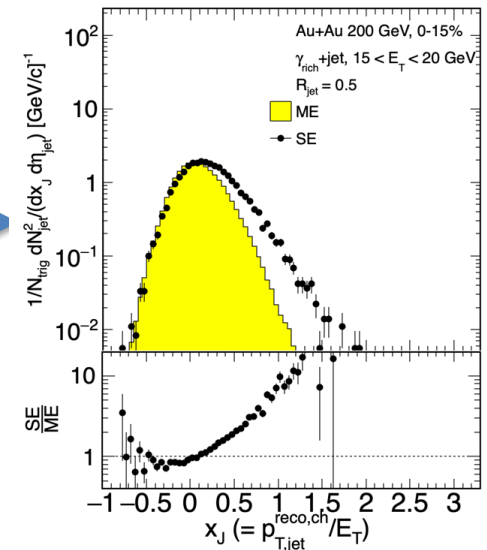
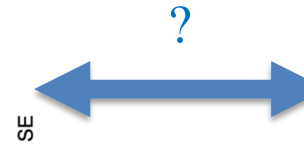
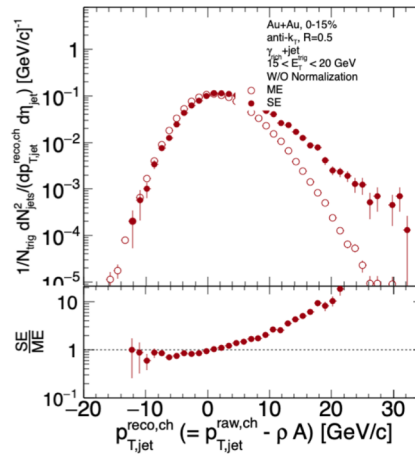
Need ME for BEMC towers

- conceptually straightforward: correct hadronic energy and then scramble
- there will no doubt be issues in practice



# What do we learn from $x_J^\gamma$ ?

What's the difference?



All we have done is rescaled the horizontal axis by  $p_T^\gamma$

- physics content is unchanged
- needs exactly the same theory calculations

We could achieve the same distribution by measuring  $p_T^{\text{jet}}$  in much narrower bins of  $p_T^\gamma$

- but will sacrifice measurement precision for unfolding etc.

Essential difference is technical: resolution in scaling by  $p_T^\gamma$

- Can  $x_J^\gamma$  be derived from the  $p_T^{\text{jet}}$  distribution using the shape of the inclusive  $\gamma$  spectrum?

$x_J^\gamma$  certainly worth measuring but benefit is incremental